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Magnetospheric Response to a Pressure Pulse in a Three-dimensional Hybrid-Vlasov Simulation

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Vlasiator Simulations

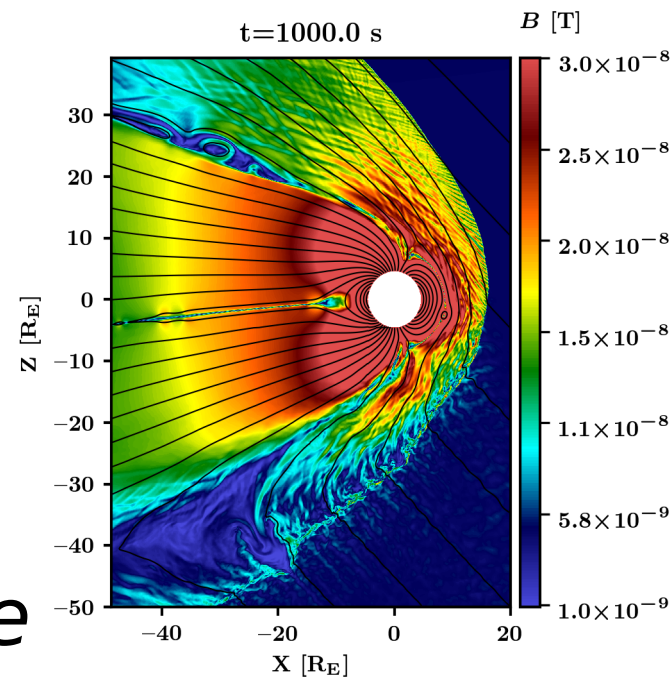
- hybrid-Vlasov (kinetic p+, fluid e-)
- 3D box (side length $\sim 100 R_E$)
- Inner boundary: $r=4.7 R_E$
- Adaptive mesh

For details:

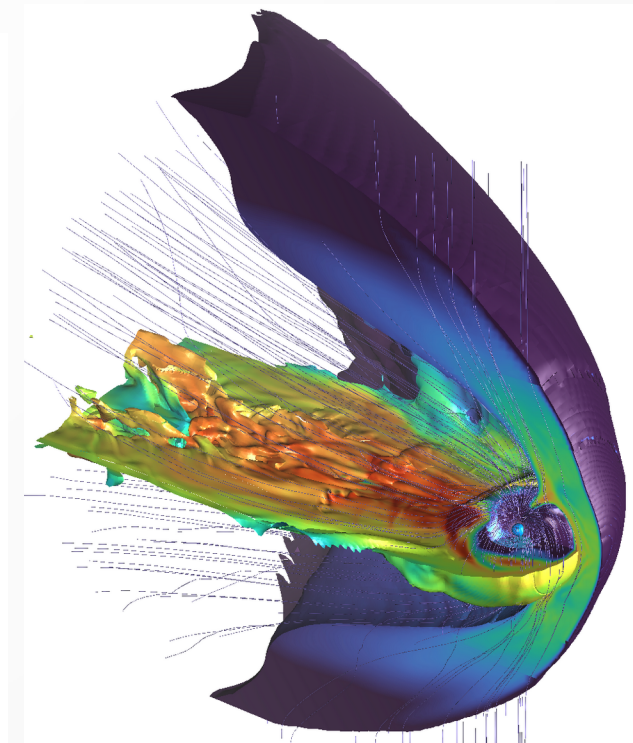
M. Palmroth et al., 2018

“Vlasov methods in space physics and astrophysics”

2D3V

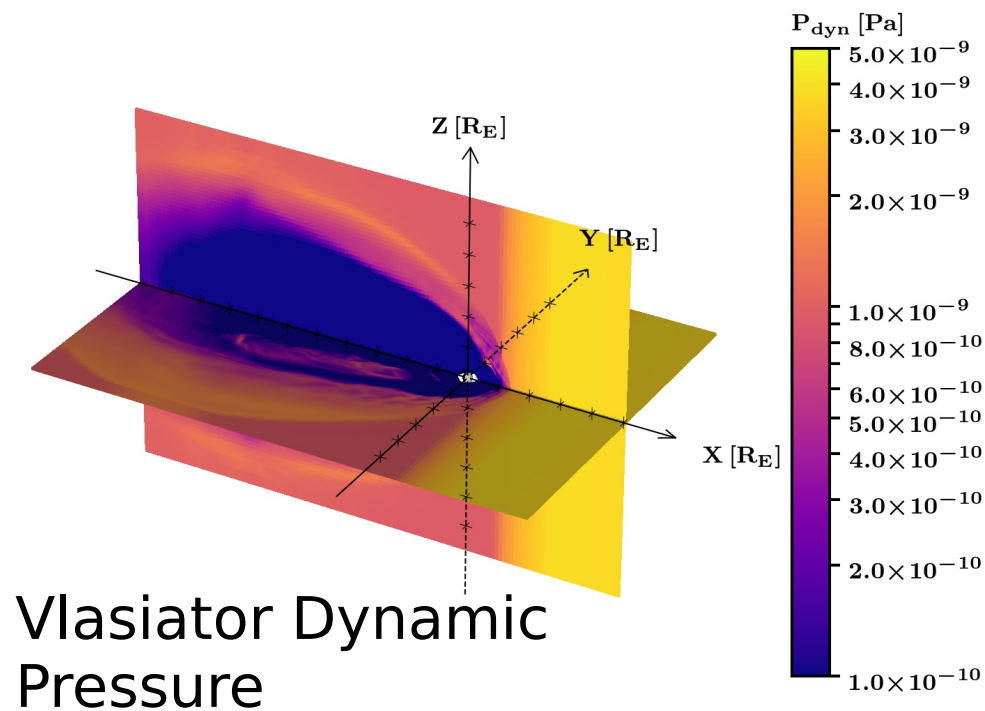


3D3V



Pressure Pulse

$t=857.0$ s – origin at (0, 0, 0) [R_E]
Tick every 10 R_E

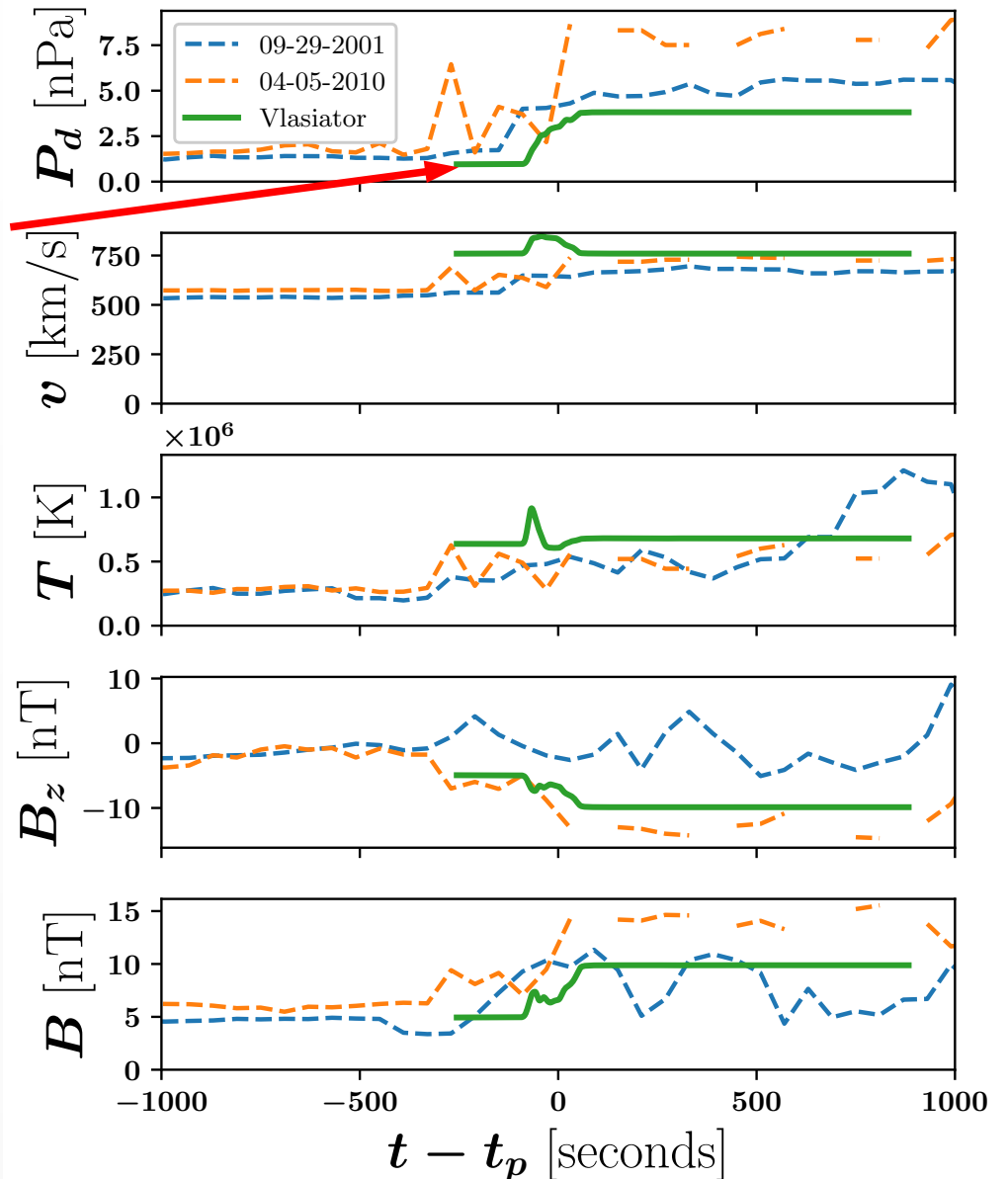


Pressure is **step function** at boundary, but **smooths out** in transit to Earth.

Not a shock (v_{sw} = constant, M_A = constant)

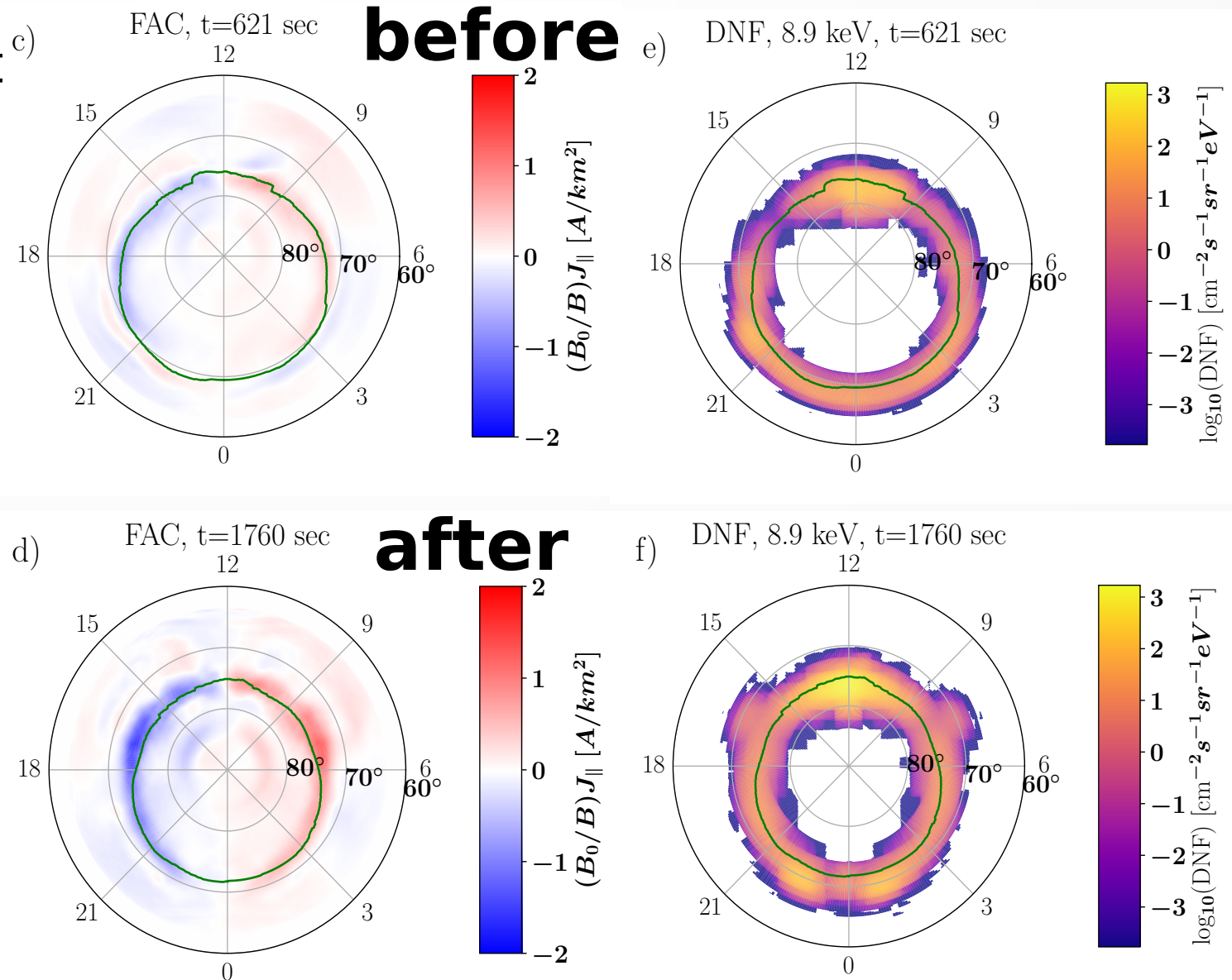
Driving	B [nT]	v_{sw} [km/s]	n [cm ⁻³]
Initial	[0, 0, -5]	750	1
Pulse	[0, 0, -10]	750	4

Virtual spacecraft near bow shock, compared with OMNI pulse events.



FACs and precipitation

- Vlasiator simulations exhibit **field-aligned currents (FACs)** and **proton precipitation** at proton aurora energies (~ 10 keV).
- Both signatures are enhanced minutes after the pressure pulse arrival, and are comparable to observations.
- Region 2 FACs not seen in this run because of simplified ionosphere boundary conditions.



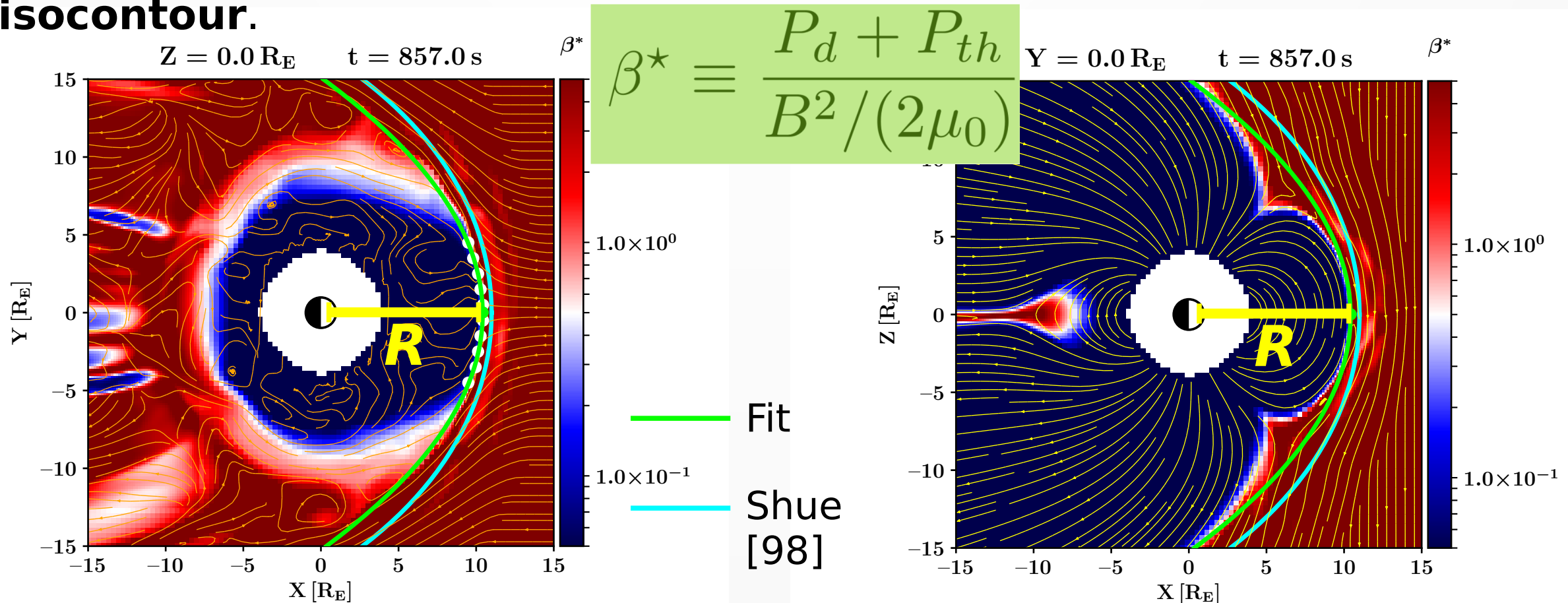
Magnetopause Identification

- Magnetopause is identified as an **isocontour** of the β^* parameter (A. Brenner et al. 2021).
- Subsolar magnetopause standoff distance R is found by fitting a parabola to the $\beta^* = 0.5$ **isocontour**.

P_d : dynamic pressure

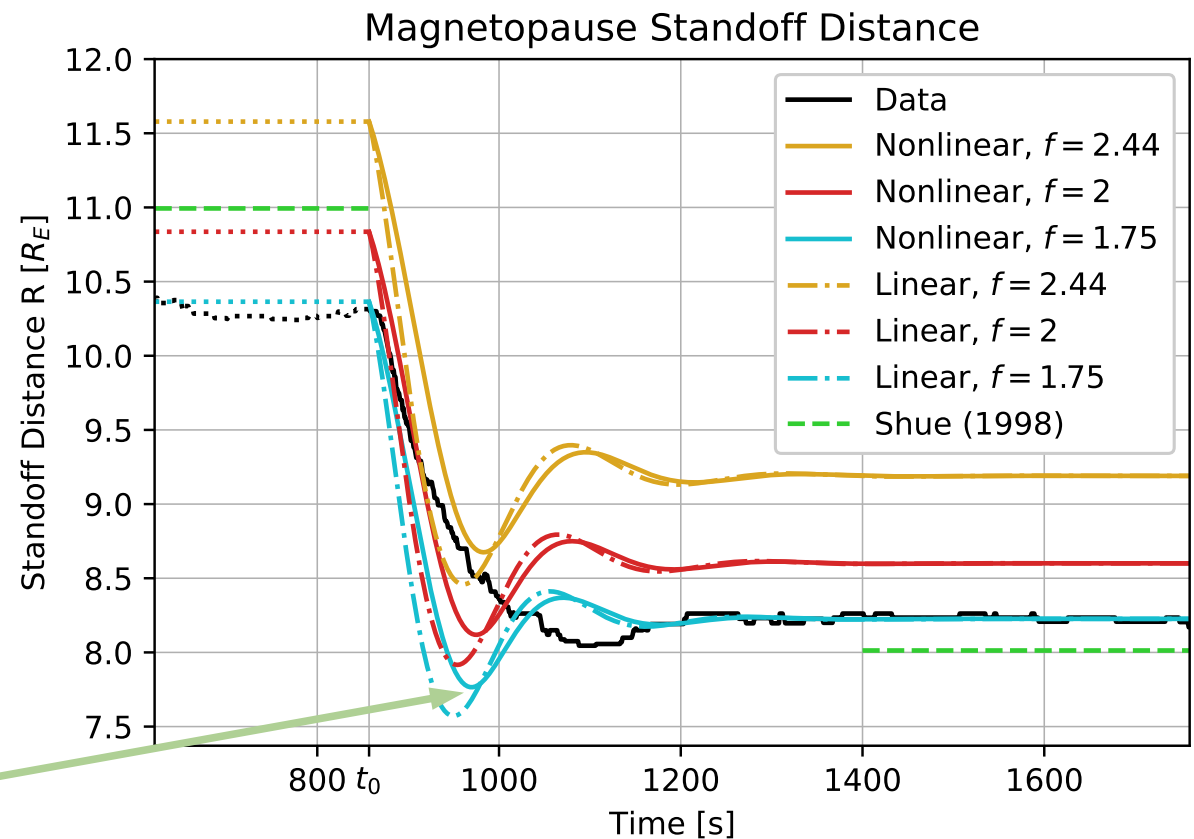
P_{th} : thermal pressure

B : Magnetic field

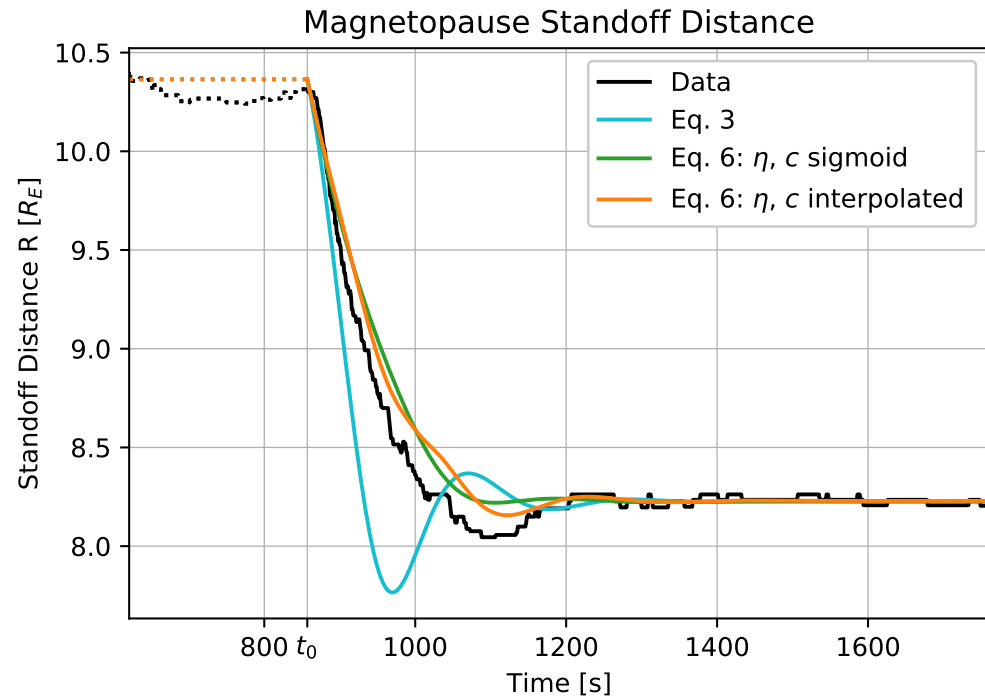
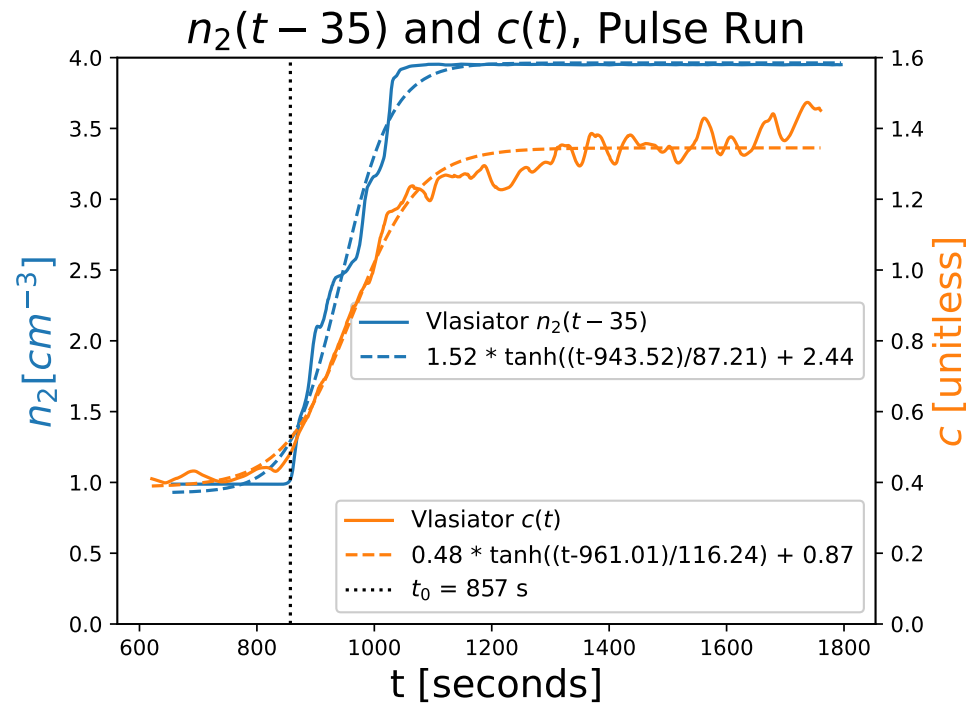


Magnetopause Oscillations

- The subsolar magnetopause standoff distance R is modeled following **Freeman et al., 1998**.
- $R(t)$ oscillates around an equilibrium set by incoming SW dynamic pressure and magnetic pressure.
- Dipole compression factor $f=1.75$ gives the best match to pre- and post-pulse equilibria.
- Both nonlinear and linearized solutions $R(t)$ reaches the global minimum value **too early**.



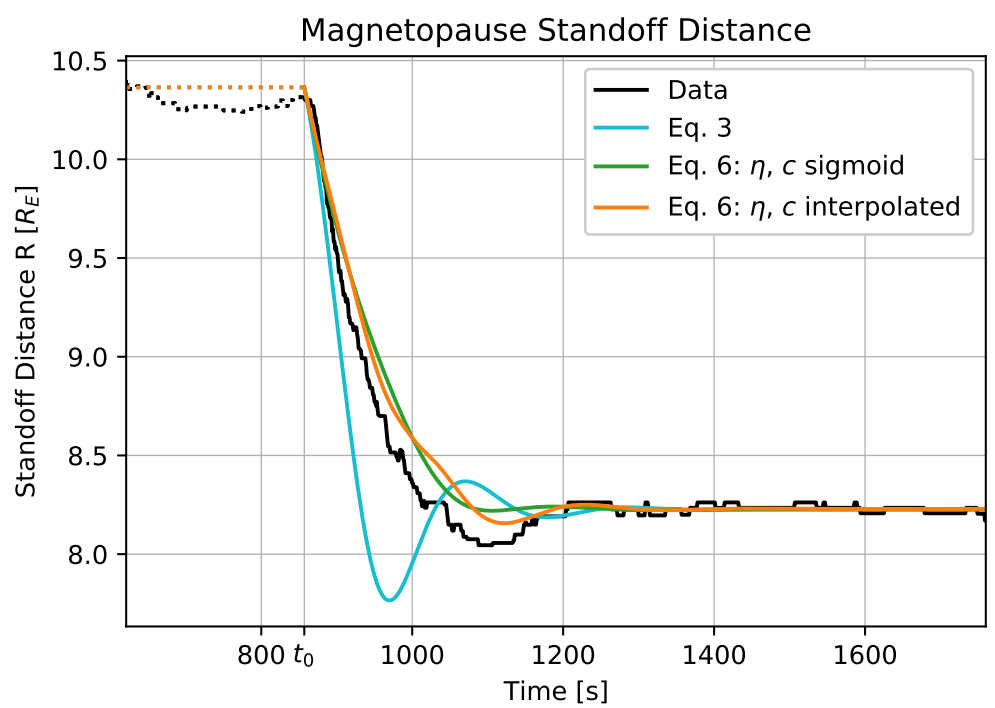
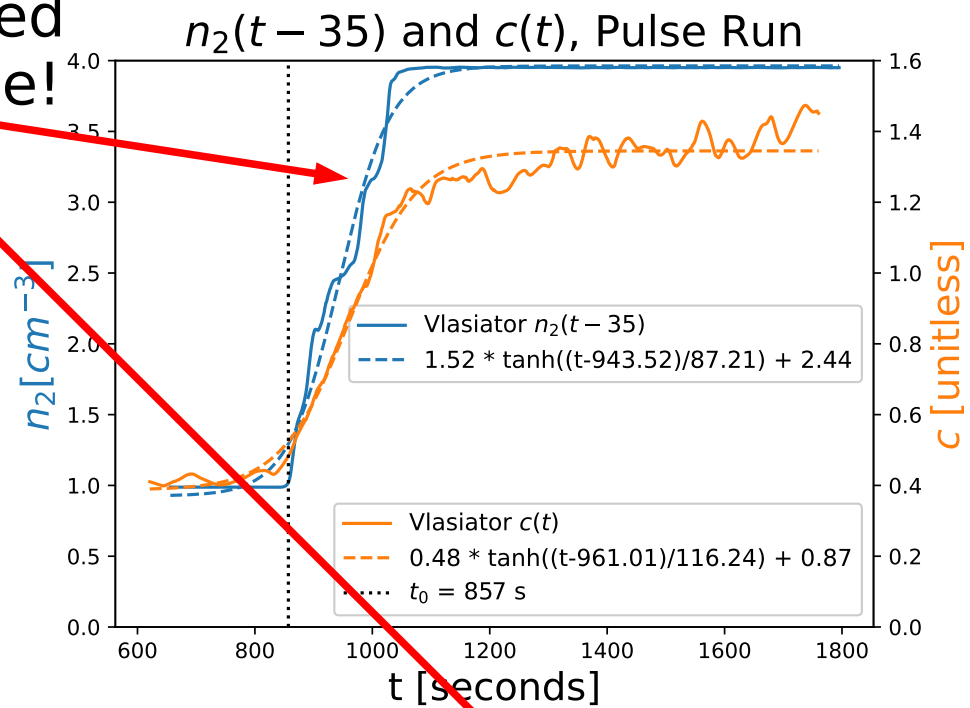
$$\frac{d^2 R}{dt^2} + \frac{1}{cR_F} \left\{ \left(v_F + \frac{dR}{dt} \right)^2 - v_F^2 \left(\frac{R_F}{R(t)} \right)^6 \right\} = 0$$



$$\frac{d^2 R}{dt^2} + \frac{1}{c(t) R_F} \left\{ \eta(t) \left(v_F + \frac{dR}{dt} \right)^2 - v_F^2 \left(\frac{R_F}{R(t)} \right)^6 \right\} = 0$$

- **GENERALIZE** Freeman et al. (1998): allow time-dependent mass loading $\mathbf{c(t)}$ and solar wind density $\mathbf{\eta(t)}$, evaluated directly from Pulse run.

Smoothed out pulse!



$$\frac{d^2 R}{dt^2} + \frac{1}{c(t) R_F} \left\{ \eta(t) \left(v_F + \frac{dR}{dt} \right)^2 - v_F^2 \left(\frac{R_F}{R(t)} \right)^6 \right\} = 0$$

- Smoothed out pulse leads magnetopause standoff **R(t)** to decrease more slowly.

Validation: Time-dependence

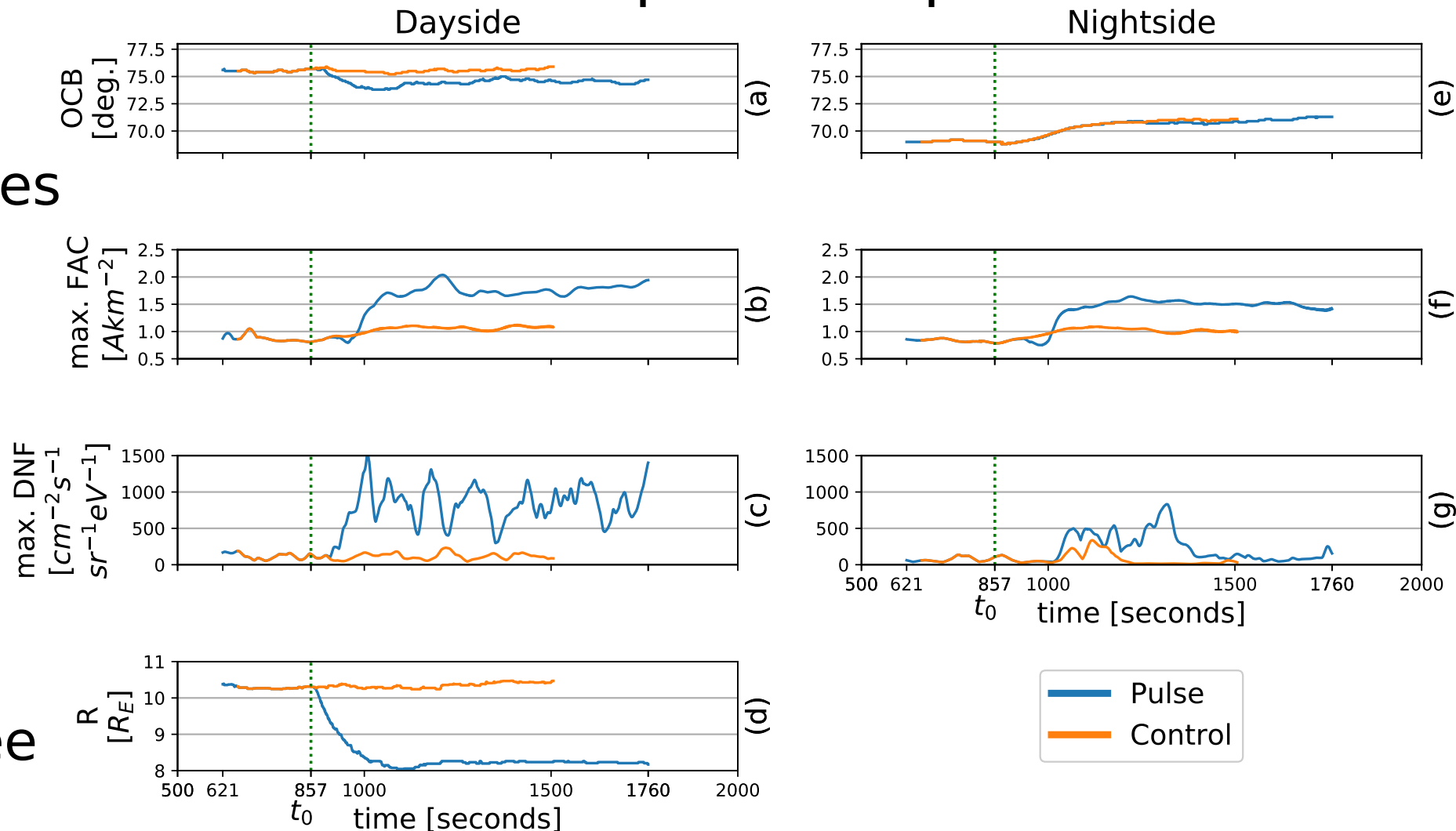
Comparison between the “Pulse” and “Control” runs allows us to isolate effects due to pressure pulse.

Only
dayside
OCB moves

FAC x2

10 keV DNF x8

Magnetopause
compression (see
previous slides)



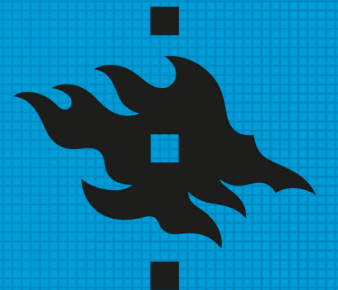
- **Vlasiator's 3D hybrid-kinetic model** of the global magnetosphere **produces expected behavior** for a pressure pulse arriving at Earth.
- The **finite transition time** of the pressure pulse causes **magnetopause oscillations to be weak and elongated** relative to established models.
- The **magnetopause oscillations are explained with a generalized model** that accounts for the finite transition time of the pressure pulse.

Horaites et al. (JGR, under review)

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